

# **Computational Fluid Dynamics**

# Lecture 5

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## **Numerical Discretization**

- Aims
  - The aim of this chapter is to identify and understand the common methods used in process of discretization
- Expected Outcomes: At the end of this chapter, students should be able to understand
  - types of numerical discretization techniques
  - numerical solutions to algebraic equations
  - How to apply common discretization techniques to typical flow equations

#### References

- 1) J. Tu, G.H. Yeoh, C. Liu, Computational Fluid Dynamics : A Practical Approach, Elsevier, 1st Edition, 2013.
- 2) C.T. Shaw, Using Computational Fluid Dynamics, Prentice Hall, 1992



### Contents

The Finite Volume Discretization Method

- Advantages and disadvantages of FVM
- Integral approximation
- Location of the solution in the finite-volume cell





# The Finite-Volume Discretization Technique

- The finite-volume method (FVM) discretized the integral form of the governing equations directly in the physical space.
- In this case, the computational domain is partitioned into a finite number of cell.
- The flow values are calculated at the centroid of each finite-volume cells.





#### Disadvantages of this method:

The FVM can approximate higher order differencing (i.e., higher than the 2<sup>nd</sup> order), however it is difficult to develop in three dimension.

#### Advantages of this method:

- The grids can be formed by combining triangular and quadrilaterals meshes for two dimensional or tetrahedral and hexahedral in three dimension
- Unstructured grids can be used or complex geometries due to its greater flexibility.
- Unlike the FDM, it doesn't require the equations to be transformed in terms of body-fitted coordinate system.





#### Finite-Volume Cell:

- A cell contains a finite (positive) volume.
- Integral equation will be approximated on the cell to form an algebraic relation.
- Size of cell indicates the level of computational resolution of the CFD analysis.
- Control surface is *faceted* into a finite number of faces.
- Face is bounded by edges.
- Edges are usually straight lines.









Assume flow variable  $\psi$  is uniform within the finite-volume shown below.

Thus:

$$\int_{V} \psi \, dV \approx V \, \psi$$

Where *V* is the volume of the finite-volume cell.



Cell *i* (*i* is an index for the cell)





> Applying Gauss's divergence theorem to a volume integral, a first order derivative of any flow variable  $\Psi$ , the differential equation along the x direction can be represented as:

$$\left(\frac{\partial \psi}{\partial x}\right) = \frac{1}{\Delta V} \int_{V} \frac{\partial \psi}{\partial x} dV = \frac{1}{\Delta V} \int_{A} \psi dA^{x} \approx \frac{1}{\Delta V} \sum_{i=1}^{N} \psi_{i} A^{x}_{i}$$

where;

- $\succ \Psi_i$  = flow values at the finite-volume cell surfaces
- $\succ$  N = number of enclosing surfaces on the finite-volume cell
- > A = surface area in the finite-volume cell.





- > For a structured grid in a 2D quadrilateral element, there are 4 enclosing surfaces (N=4) of the finitevolume cell. On the other hand, in 3D, N = 6.
- > Similar to the x-direction , the first-order derivative for  $\Psi$  in the y-direction will have the form:

$$\left(\frac{\partial \psi}{\partial y}\right) = \frac{1}{\Delta V} \int_{\Delta V} \frac{\partial \psi}{\partial y} dV = \frac{1}{\Delta V} \int_{A} \psi dA^{y} \approx \frac{1}{\Delta V} \sum_{i=1}^{N} \psi_{i} A^{y}_{i}$$





Position of the flow and heat transfer solution in the finitevolume cell with respect to the mesh may have two forms:

i. Node-Centered: The flow and heat transfer solutions are positioned at the vertices of the mesh grid. The finite-volume cell is formed about the vertex.







ii. Cell-Centered: The flow and heat transfer solutions are positioned at the centroid of the cell





#### Example



Determine the discreditised form of the following two dimensional continuity equation using the FVM.

 $\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0$ Use structured uniform grid arrangement.





### **Solution**



Applying the finite volume equations and either of the solution locations (node-centered or cell-centered) provided in the previous slides, the following expressions can be obtained

$$\frac{\partial u}{\partial x} \approx \frac{1}{\Delta V} \sum_{i=1}^{4} u_i A_i^x$$
$$= \frac{1}{\Delta V} \left( u_e A_e^x - u_w A_w^x + u_n A_n^{x=0} - u_s A_s^{x=0} \right)$$

$$\frac{\partial v}{\partial y} \approx \frac{1}{\Delta V} \sum_{i=1}^{4} v_i A_i^y$$
$$= \frac{1}{\Delta V} \left( v_e A_e^{y=0} - v_w A_w^{y=0} + v_n A_n^y + v_s A_s^y \right)$$





- > The face velocities  $\mathcal{U}_{e}$ ,  $\mathcal{U}_{w}$ ,  $\mathcal{V}_{n}$  and  $\mathcal{V}_{s}$  are positioned halfway between centroid of each finite-volume cells,
- This makes it easy for us to calculate the face velocities from the values located at the centroid of each finite-volume cells.

➤ Thus:

$$u_{e} = \frac{u_{P} + u_{E}}{2}; \quad u_{w} = \frac{u_{P} + u_{W}}{2}; \quad v_{n} = \frac{v_{P} + v_{N}}{2}; \quad v_{s} = \frac{v_{P} + v_{S}}{2}$$

By substituting the above expressions to the discretized form of the velocity first-order derivatives, the final form of the discretized continuity equation becomes:

$$\left(\frac{\mathcal{U}_{P}+\mathcal{U}_{E}}{2}\right)A_{e}^{x}-\left(\frac{\mathcal{U}_{P}+\mathcal{U}_{W}}{2}\right)A_{w}^{x}+\left(\frac{\mathcal{V}_{P}+\mathcal{V}_{N}}{2}\right)A_{n}^{y}-\left(\frac{\mathcal{V}_{P}+\mathcal{V}_{S}}{2}\right)A_{s}^{y}=0$$





>  $A_e^x = A_w^x = \Delta y$  and  $A_n^y = A_s^y = \Delta x$ , the above equation can then be expressed by:

$$\left(\frac{\mathcal{U}_{P}+\mathcal{U}_{E}}{2}\right)A_{e}^{x}-\left(\frac{\mathcal{U}_{P}+\mathcal{U}_{W}}{2}\right)A_{w}^{x}+\left(\frac{\mathcal{V}_{P}+\mathcal{V}_{N}}{2}\right)A_{n}^{y}-\left(\frac{\mathcal{V}_{P}+\mathcal{V}_{S}}{2}\right)A_{s}^{y}=0$$

 $\succ$  and reduced to:

$$\left(\frac{\mathcal{U}_{P}+\mathcal{U}_{E}}{2}\right)\Delta y - \left(\frac{\mathcal{U}_{P}+\mathcal{U}_{W}}{2}\right)\Delta y + \left(\frac{\mathcal{V}_{P}+\mathcal{V}_{N}}{2}\right)\Delta x - \left(\frac{\mathcal{V}_{P}+\mathcal{V}_{s}}{2}\right)\Delta x = 0$$

> or in another form:

$$\left(\frac{\boldsymbol{\mathcal{U}}_E - \boldsymbol{\mathcal{U}}_W}{2}\right) \Delta y + \left(\frac{\boldsymbol{\mathcal{V}}_N - \boldsymbol{\mathcal{V}}_S}{2}\right) \Delta x = 0$$

$$\frac{\boldsymbol{\mathcal{U}}_E - \boldsymbol{\mathcal{U}}_W}{2\Delta x} + \frac{\boldsymbol{\mathcal{V}}_N - \boldsymbol{\mathcal{V}}_S}{2\Delta y} = 0$$



## Summary

- FVM can be applied to greater than the second order representation of the flow.
- $\succ$  First order derivative of any flow variable  $\psi$

$$\left(\frac{\partial \psi}{\partial x}\right) = \frac{1}{\Delta V} \int_{V} \frac{\partial \psi}{\partial x} dV = \frac{1}{\Delta V} \int_{A} \psi dA^{x} \approx \frac{1}{\Delta V} \sum_{i=1}^{N} \psi_{i} A^{x}_{i}$$

$$\left(\frac{\partial \psi}{\partial y}\right) = \frac{1}{\Delta V} \int_{\Delta V} \frac{\partial \psi}{\partial y} dV = \frac{1}{\Delta V} \int_{A} \psi dA^{y} \approx \frac{1}{\Delta V} \sum_{i=1}^{N} \psi_{i} A^{y}_{i}$$

Location of cells

- Node-Centered
- Cell-Centered







### Dr. A. Nurye Research interest:

- Computational Fluid Dynamics,
- Thermo-fluids,
- Multidisciplinary Numerical Modelling and Simulation

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