

# Computational Fluid Dynamics

## Lecture 3

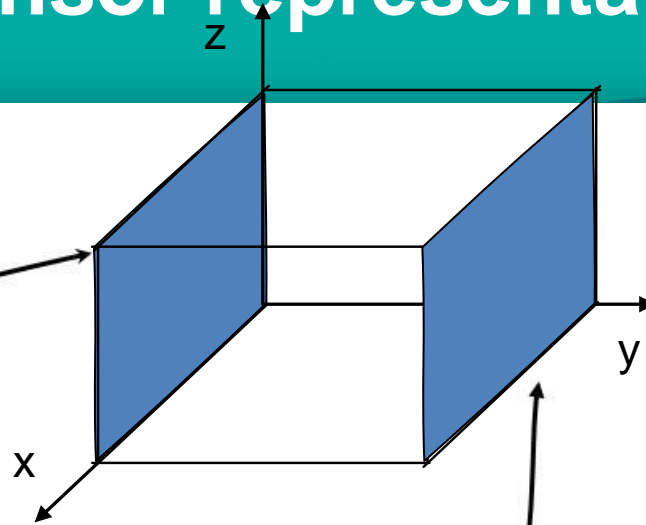
by

Dr. A. Nurye  
Faculty of Mechanical Engineering  
[nurye@ump.edu.my](mailto:nurye@ump.edu.my)

# Conservation of Momentum and Energy

- Aims
  - The aim of this lecture is to recall the basic momentum equations related to fluid flow and heat transfer.
- Expected Outcomes: At the end of this lecture, students should be able to
  - understand the governing momentum equations
  - derive 2D momentum equations
- References
  - 1) J. Tu, G.H. Yeoh, C. Liu, Computational Fluid Dynamics : A Practical Approach, Elsevier, 1st Edition, 2013.
  - 2) C.T. Shaw, Using Computational Fluid Dynamics, Prentice Hall, 1992

# Tensor representation



Each of the six faces has a **direction**.

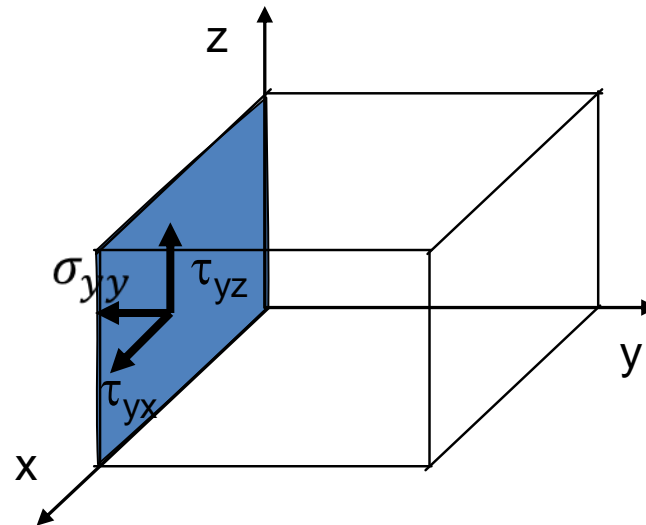
For example, this face

and this face

are normal to the **y direction**

A force acting on any face can act in the x, y and z directions.

Consider the face below.



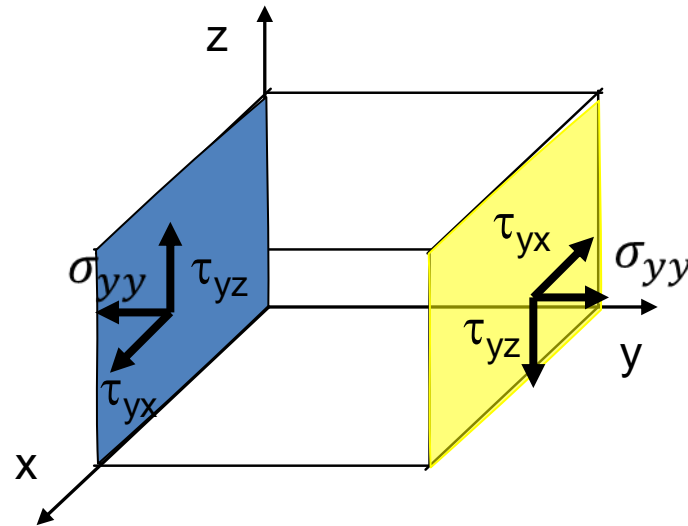
The face is perpendicular to the  $y$  direction.

The force per unit face area acting in the  $x$  direction on that face is the stress  $\tau_{yx}$  (first component: face, second component: stress direction).

The forces per unit face area acting in the  $y$  and  $z$  directions on that face are the stresses  $\tau_{yy}$  and  $\tau_{yz}$ .

Here  $\tau_{yy}$  is a **normal stress** (acts normal, or perpendicular to the face) and  $\tau_{yx}$  and  $\tau_{yz}$  are **shear stresses** (act parallel to the face)

Some conventions are in order



Normal stresses are defined to be positive **outward**, so the orientation is reversed on the face located  $\Delta y$  from the origin

Shear stresses similarly reverse sign on the opposite face.

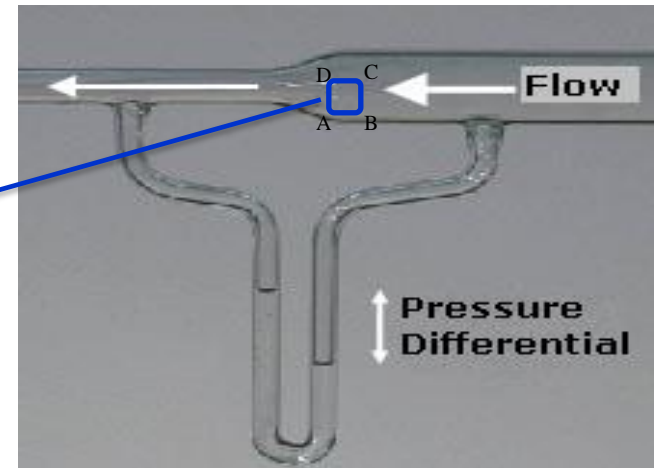
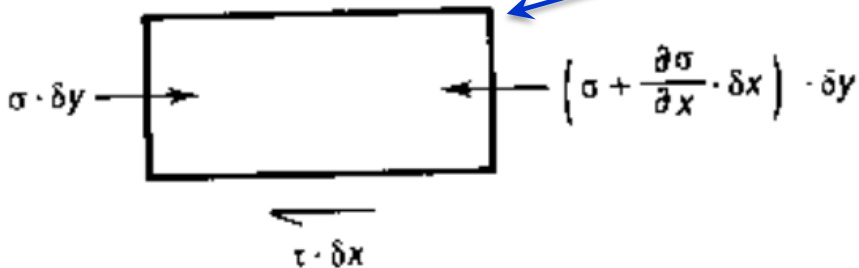
Thus a positive normal stress puts a body in tension, and a negative normal stress puts the body in compression. Shear stresses always put the body in shear.

## b) Conservation of momentum

- The basis for the derivation of momentum equations the Newton's 2<sup>nd</sup> Law of Motion
- It is function of the forces on the applied to the fluid and the its acceleration.

Rate of increase of momentum of fluid particle = Sum of forces on fluid particle

$$\tau \cdot \delta x + \frac{\partial \tau}{\partial y} \cdot \delta y \cdot \delta x$$



<https://commons.wikimedia.org/wiki/File:VenturiFlow.png>

Consider small portion of the domain fluid as shown above, where  $\tau$  is the shear stress (parallel to the surfaces) and  $\sigma$  is normal stress

- The acceleration of a fluid particle can be calculated by taking change of velocity with spatial directions (  $x$  and  $y$  ), and time (  $t$  ).
- Let  $u$  and  $v$  be the  $x$  and  $y$ , respectively, components of the fluid velocity.
- Thus, the material derivative of  $u$  will have the form

$$Du = \frac{\partial u}{\partial x} dx + \frac{\partial u}{\partial y} dy + \frac{\partial u}{\partial t} dt$$

- which becomes, on dividing by delta  $t$

$$\frac{du}{dt} = \frac{\partial u}{\partial x} \frac{dx}{dt} + \frac{\partial u}{\partial y} \frac{dy}{dt} + \frac{\partial u}{\partial t} \frac{dt}{dt}$$

➤ Now,  $\frac{dx}{dt} = u$  , and,  $\frac{dy}{dt} = v$

➤ Thus,  $\frac{Du}{Dt} = u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + \frac{\partial u}{\partial t}$

➤ This expression represents the total acceleration of the fluid in the  $x$ -direction and is called the material derivative of  $u$  .

➤ Newton's Second Law: 
$$F_x = ma_x = m \frac{Du}{Dt} = m \left( u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + \frac{\partial u}{\partial t} \right)$$

where  $F_x$ ,  $m$  and  $a_x$  are the force, mass, and acceleration of the fluid in the  $x$ -direction.



- The x-component of the force applied on the fluid element can thus be calculated as

$$\rho \left( u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + \frac{\partial u}{\partial t} \right) dx dy = \frac{\partial \sigma}{\partial x} dx dy + \frac{\partial \tau}{\partial y} dx dy$$

- But

$$\sigma = -p + 2\mu \frac{\partial u}{\partial x}, \quad \text{and} \quad \tau = \left( \frac{\partial v}{\partial x} + \frac{\partial u}{\partial y} \right)$$

- Substituting the normal and shear stresses we get the x-momentum equation

$$\rho \left( \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} \right) = -\frac{\partial p}{\partial x} + \mu \left( \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right)$$

- Similarly, the y-momentum equation is

$$\rho \left( \frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} \right) = -\frac{\partial p}{\partial y} + \mu \left( \frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} \right)$$

Where  $\mu$  is the viscosity of the fluid and  $\rho$  is its density.

## Example

Consider a 2D steady, incompressible viscous flow between two parallel plates a distance  $L$  apart. Assume that the plates are very wide and very long, so that the flow between the plates is essentially in the  $x$  direction. The upper plate moves at velocity  $V$ . Neglecting the pressure and gravity effect, determine

- a) the velocity field,
- b) the shear force per unit area acting on the bottom plate,

# Solution

- Step 1: List the given information
  - Assumptions
    1. The plates are infinite in  $x$  and  $z$  directions
    2. Steady state flow condition,  $\partial/\partial t = 0$
    3.  $v=w=0$ , the flow is assumed to be parallel
    4. Incompressible flow
    5. No pressure gradient in  $y$ -direction
    6. Two dimensional flow,  $w=0$ ,  $\partial/\partial z = 0$
    7. No Gravity ,
  
  - Boundary conditions (Take the origin at the corner of the bottom plate)
    1. At  $y=0$  :  $u=0$ ,  $v=0$ ,  $w=0$
    2. At  $y=h$  :  $u=V$ ,  $v=0$ ,  $w=0$

➤ Step 2: Simplify

Continuity

$$\frac{\partial U}{\partial x} + \cancel{\frac{\partial V}{\partial y}} + \cancel{\frac{\partial W}{\partial z}} = 0$$

$$\boxed{\frac{\partial U}{\partial x} = 0}$$

X-momentum

$$\rho \left( \cancel{\frac{\partial U}{\partial t}} + U \cancel{\frac{\partial U}{\partial x}} + V \cancel{\frac{\partial U}{\partial y}} + W \cancel{\frac{\partial U}{\partial z}} \right) = -\cancel{\frac{\partial P}{\partial x}} + \cancel{\rho g_x} + \mu \left( \cancel{\frac{\partial^2 U}{\partial x^2}} + \frac{\partial^2 U}{\partial y^2} + \cancel{\frac{\partial^2 U}{\partial z^2}} \right)$$

$$\boxed{\frac{\partial^2 U}{\partial x^2} = 0}$$

## ➤ Step 3: Integrate

X-momentum

$$\frac{d^2 u}{dy^2} = 0 \xrightarrow{\text{integrate}} \frac{du}{dy} = C_1 \xrightarrow{\text{integrate}} u(y) = C_1 y + C_2$$

## ➤ Step 4: Apply the BC's

$$\begin{aligned} > y=0, u=0, & \quad C_1(0) + C_2 \Rightarrow \underline{C_2 = 0} \\ > y=h, u=V, & \quad C_1 h \Rightarrow \underline{C_1 = V/h} \end{aligned}$$

This gives

$$u(y) = V \frac{y}{h}$$

- Step 5: Substitute the velocity values in to the continuity and momentum equations and verify that both equations are conserved.
- From the x, y, and z, velocity components

$$u = \frac{vy}{h}, \quad v = 0, \quad \text{and} \quad w = 0$$

- We get  $\frac{\partial u}{\partial x} = 0, \frac{\partial v}{\partial y} = 0, \frac{\partial w}{\partial z} = 0$

- Continuity is satisfied

$$0 + 0 + 0 = 0$$

- X-momentum is satisfied

$$\rho \left( 0 + V \frac{y}{h} \cdot 0 + 0 \cdot V/h + 0 \cdot 0 \right) = -0 + \rho \cdot 0 + \mu (0 + 0 + 0)$$

$$0 = 0$$

- Finally, the shear force on bottom plate can be calculated as

$$\tau_{ij} = \begin{pmatrix} 2\mu \frac{\partial U}{\partial x} & \mu \left( \frac{\partial U}{\partial y} + \frac{\partial V}{\partial x} \right) & \mu \left( \frac{\partial U}{\partial z} + \frac{\partial W}{\partial x} \right) \\ \mu \left( \frac{\partial V}{\partial x} + \frac{\partial U}{\partial y} \right) & 2\mu \frac{\partial V}{\partial y} & \mu \left( \frac{\partial V}{\partial z} + \frac{\partial W}{\partial y} \right) \\ \mu \left( \frac{\partial W}{\partial x} + \frac{\partial U}{\partial z} \right) & \mu \left( \frac{\partial W}{\partial y} + \frac{\partial V}{\partial z} \right) & 2\mu \frac{\partial W}{\partial z} \end{pmatrix} = \begin{pmatrix} 0 & \mu \frac{V}{h} & 0 \\ \mu \frac{V}{h} & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$\frac{\vec{F}}{A} = \tau_w = \mu \frac{V}{h} \hat{i}$$



# CFD EXERCISE

Perform a 2D CFD analysis for the laminar incompressible flow of air between two parallel plates discussed in the previous example. The length of the plates is 2 m and the gap between them is 0.2 m. The air is flowing in the horizontal direction along the plate length at velocity of 0.025 m/s.

# Conclusion of the lecture

- The x and y components of a two dimensional conservation of momentum equation

$$\rho \left( \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} \right) = -\frac{\partial p}{\partial x} + \mu \left( \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right)$$

$$\rho \left( \frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} \right) = -\frac{\partial p}{\partial y} + \mu \left( \frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} \right)$$



# Dr. A. Nurye

## Research interest:

- Computational Fluid Dynamics,
- Thermo-fluids,
- Multidisciplinary Numerical Modelling and Simulation

## Contact:

Tel: +094246259

email: [nurye@ump.edu.my](mailto:nurye@ump.edu.my)